

*Article*

# **Composites Based on Electrodeposited WO<sup>3</sup> and TiO<sup>2</sup> Nanoparticles for Photoelectrochemical Water Splitting**

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 $A$ bstract: Photoelectrochemically active  $\mathsf{WO}_3$  films were fabricated by electrodeposition from an acidic (pH 2), hydrogen-peroxide-containing electrolyte at  $-0.5$  V vs. SCE. WO<sub>3</sub>-TiO<sub>2</sub> composites were then synthesized under the same conditions, but with 0.2  $g/L$  of anatase TiO<sub>2</sub> nanoparticles ( $\oslash$  36 nm), mechanically suspended in the solution by stirring. After synthesis, the films were annealed at 400 ◦C. Structural characterization by XRD showed that the WO<sub>3</sub> films exhibit the crystalline structure of a non-stoichiometric hydrate, whereas, in  $\rm WO_3$ -Ti $\rm O_2$ , the  $\rm WO_3$  phase was monoclinic. The oxidation of tungsten, as revealed by XPS, was  $W^{6+}$  for both materials. Ti was found to exist mainly as  $Ti^{4+}$ in the composite, with a weak  $Ti^{3+}$  signal. The efficiency of the WO<sub>3</sub> films and composites as an oxygen evolution reaction (OER) photo-electrocatalyst was examined. The composite would generate approximately three times larger steady-state photocurrents at 1.2 V vs. SCE in a neutral  $0.5 M Na<sub>2</sub>SO<sub>4</sub>$ electrolyte compared to  $WO<sub>3</sub>$  alone. The surface recombination of photogenerated electron–hole pairs was characterized by intensity-modulated photocurrent spectroscopy (IMPS). Photogenerated charge transfer efficiencies were calculated from the spectra, and at 1.2 V vs. SCE, were 86.6% for WO<sub>3</sub> and 62% for WO<sub>3</sub>-TiO<sub>2</sub>. Therefore, the composite films suffered from relatively more surface recombination but generated larger photocurrents, which resulted in overall improved photoactivity.

**Keywords:** tungsten trioxide titania composite; electrodeposition; photoelectrochemical water splitting; IMPS of OER

# **1. Introduction**

The adverse environmental impact of using fossil fuels for energy needs and generating  $H_2$  for chemical processing has led to significant research toward finding alternative and sustainable energy carriers, such as hydrogen  $(H<sub>2</sub>)$ , with  $H<sub>2</sub>$  produced by solar water splitting is more advantageous compared to other forms of green hydrogen  $[1-3]$  $[1-3]$ . The photoelectrochemical (PEC) water-splitting process includes light harvesting, photogenerated charge separation, and surface catalysis, and is an effective strategy as an alternative to fossil fuels. The reduction of water to produce hydrogen is coupled with the more kinetically hindered oxygen evolution reaction, which is typically carried out on metal oxide photoanodes  $[4,5]$  $[4,5]$ . Although TiO<sub>2</sub> is recognized as an effective inexpensive photocatalyst in the UV range,  $[6,7]$  $[6,7]$ , the addition of tungsten trioxide (WO<sub>3</sub>) is widely known to lower the band gap  $(\sim 2.75 \text{ eV})$  and hence, allows absorption to occur within the visible range of the solar spectrum,  $[8-11]$  $[8-11]$ . However, WO<sub>3</sub> alone suffers from high recombination of photon-generated electron–hole pairs [\[12\]](#page-12-8) and is susceptible to photocorrosion caused by peroxo-species formed on the surface during the photoelectrochemical reaction [\[13\]](#page-12-9). Tungsten trioxide doped by TiO<sub>2</sub> shows a red-shifted optical response in the visible light



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range due to a band gap decrease from a lowering of the conduction band edge, and can act as a photocatalyst. Momeni and Nazari [\[14\]](#page-12-10) reported a change in band gap from 3.23 eV for bare TiO<sub>2</sub> to 2.78 eV for WO<sub>3</sub>-doped TiO<sub>2</sub> nanotubes. Also, doping TiO<sub>2</sub> with cations like tungsten ( $W^{6+}$ ) introduces energy levels in TiO<sub>2</sub> just below the conduction band that acts as an electron trap increasing the lifetime of charge carriers, and thus retarding recombination [\[15\]](#page-12-11).

Heterojunctions of  $WO_3/TiO_2$  layers have been used to increase the efficiency of the photocatalysts [\[16](#page-12-12)[–20\]](#page-12-13) by aligning the band structures of the two coupled semiconductors to encourage migration of photogenerated electrons (e−) and holes (h<sup>+</sup> ) in separate directions across the heterojunction boundary, which reduces the electron–hole recombination. The highly oxidative valence band of  $WO_3$ , arising from the O 2p orbitals, also helps to drive oxidation reactions, such as PEC oxygen evolution or degradation of dyes.  $WO<sub>3</sub>/TiO<sub>2</sub>$  composites have been reported to perform as stable and efficient photoanodes in water splitting [\[21–](#page-12-14)[23\]](#page-12-15). Numerous reports of dye degradation with tungsten doped  $TiO<sub>2</sub>$  or heterojunctions are reported through the reactions of a photogenerated superoxide [\[24–](#page-13-0)[28\]](#page-13-1). One such heterojunction unique to electrodeposition is the work pre-sented by Martin et al. [\[29\]](#page-13-2). They electrodeposited  $WO_3$  onto a metallic Ti electrode then heat treated it to create a  $WO_3/TiO_2/Ti$  layered material that exhibited high photocurrent values toward the photoelectrocatalytic degradation of propyl paraben. An alternative strategy to reduce recombination rates is to apply a potential or bias and photogenerated electrons are drawn away from the electrode–electrolyte interface through the external surface. This electrically enhanced photocatalysis was observed for electrochemically prepared  $TiO<sub>2</sub>/WO<sub>3</sub>$  layers on stainless steel for the oxidation of organics [\[30\]](#page-13-3).

Nanostructuring is also important for generating quantum confinement effects that alter the electron and hole transport behaviors, shifting of the electronic band structure, and the large surface-to-volume ratio can significantly increase the surface reaction sites [\[31\]](#page-13-4). However, it can also introduce defects or additional grain boundaries that can promote recombination [\[32\]](#page-13-5).

As mentioned, synergistic effects, as a result of combining different photoactive materials, provide a promising pathway to improve the photoelectrochemical performance in heterojunction devices for solar water splitting [\[33\]](#page-13-6). To this end, a nanocomposite WO<sub>3</sub>-TiO<sub>2</sub> film was electrodeposited with a WO<sub>3</sub> matrix incorporating nanoparticles of  $TiO<sub>2</sub>$ , in contrast to layered materials or those with a matrix of  $TiO<sub>2</sub>$ . The electrodeposition of WO<sup>3</sup> films followed the approach by Pauporté [\[34\]](#page-13-7), where the electrochemical deposition mechanism is multistaged and occurs through the reduction of a peroxy-tungstate species as outlined in Equations (1) to (3) [\[35](#page-13-8)[–37\]](#page-13-9).

$$
2\text{WO}_4^{2-} + 4\text{H}_2\text{O}_2 \rightarrow \text{W}_2\text{O}_{11}^{2-} + 2\text{OH}^- + 3\text{H}_2\text{O}
$$
 (1)

$$
W_2O_{11}^{2-} + (2+3x)H^+ + 3xe^- \rightarrow 2H_xWO_3 + \frac{(2+x)}{2}H_2O + \frac{(8-x)}{4}O_2
$$
 (2)

$$
H_xWO_3 \to WO_3 + xH^+ + xe^-
$$
 (3)

In this study, the approach is modified by adding  $TiO<sub>2</sub>$  nanoparticles to the sodium tungstate–hydrogen peroxide electrolyte. The electrodeposition is then carried out to simultaneously electrodeposit a  $WO_3$  film and capture TiO<sub>2</sub> nanoparticles, resulting in a one-step synthesis of a composite with improved photoactivity for water splitting. Moreover, the use of an advanced non-stationary photoelectrochemical technique (IMPS) is demonstrated to probe charge carrier photogeneration in the space-charge layer and to elucidate important kinetic constants for charge transfer and recombination.

# **2. Materials and Methods**

# *2.1. Composite Electrodeposition*

A schematic representation of the cell and process used to deposit  $WO_3$  and  $WO_3$ -TiO<sub>2</sub> films is shown in Figure [1.](#page-2-0) All used chemicals were of ACS reagent grade. The electrolyte

to electrodeposit WO<sub>3</sub> films was prepared from 0.025 M Na<sub>2</sub>WO<sub>4</sub>·2H<sub>2</sub>O (Fisher Chemical, Hampton, NH, USA),  $0.075\%$  ( $v/v$ )  $\rm H_2O_2$  (Fisher Chemical, Hampton, NH, USA) which was 1.25 mL of 30% hydrogen peroxide, and the pH was adjusted to 1.4 with sulfuric acid (Fisher Cchemical, Hampton, NH, USA). To electrodeposit composite  $\rm WO_3\text{-}TiO_2$ films, the same electrolyte composition was used, but additionally, a particle loading of 0.2 g/L of anatase TiO<sub>2</sub> (Thermo Fisher Scientific, Waltham, MA, USA) of average ⊘ 36 nm was added. The nanoparticles were mechanically suspended in the electrolyte by using a magnetic stir bar. Electrodeposition was carried out in a three-electrode cell, with a stainless-steel (type 316) working electrode (plates of  $2 \text{ cm} \times 2 \text{ cm}$  dimensions were used as the substrate, with a circular working area of  $1.226 \text{ cm}^2$ ), a Pt counter electrode, and an SCE reference electrode. The substrate was positioned in a holder parallel to the bottom of the electrodeposition bath. Deposition occurred during an applied potential of  $-0.5$  V vs. SCE for 5 min.

 $\overline{\phantom{a}}$  films is shown in Figure 1. All used chemicals were of  $\overline{\phantom{a}}$ 

<span id="page-2-0"></span>

**Figure 1.** Schematic representation of material synthesis and annealing process. **Figure 1.** Schematic representation of material synthesis and annealing process.

# *2.2. Structural and Surface Morphology Characterization 2.2. Structural and Surface Morphology Characterization*

The as-deposited WO<sub>3</sub> and WO<sub>3</sub>-TiO<sub>2</sub> films were thermally annealed at 400 °C for two hours, with a 2  $\degree$ C min<sup>-1</sup> temperature ramp-up rate. The surface morphology was characterized by a Hitachi TM 3000 scanning electron microscope (Hitachi Ltd., Tokyo, Japan).<br>NGC 1166  $XRD$  diffraction patterns were obtained with a Rigaku MiniFlex II x-ray diffractometer (Rigaku Corporation, Tokyo, Japan).

tometer (Rigaku Corporation, Tokyo, Japan). The XPS analyses were carried out with a Kratos Axis Supra spectrometer (Kratos The  $\lambda_1$  S analyses were carried out with a Kratos  $A \lambda_1$ s Supra spectrometer (Kratos  $\lambda_1$ Analytical Limited, Manchester, UK), using a monochromatic Al K(alpha) source (15 mA, Analytical Limited, Manchester, UK), using a monochromatic Al K(alpha) source (15 mA,  $15 \text{ KV}$ . Survey scan analyses were carried out on the area of  $300 \times 700$   $\mu$ m at a pass energy of 160 eV. High-resolution analyses were also carried out on the area of  $300 \times 700$   $\mu$ m, but of 160 eV. High-resolution analyses were also carried out on the area of 300  $\times$  700  $\mu$ m, but on the area of 300  $\times$  700  $\times$  7 at a pass energy of 20 eV. The XPS signal due to adventitious carbon located at 284.8 eV<br>use used as a binding energy (PE) reference was used as a binding energy (BE) reference. was used as a binding energy (BE) reference. 15 kV). Survey scan analyses were carried out on the area of  $300 \times 700$  µm at a pass energy

# *2.3. PEC Characterization*

*2.3. PEC Characterization*  The photoelectrochemical response of the resulting deposits was examined with linear sweep voltammetry, chronoamperometry, and IMPS, using a ZAHNER Zennium CIMPS-QE/IPCE system (Zahner-Elektrik, Kronach, Germany). The photoelectrochemical cell included a quartz window for front-illumination of the films, a platinum counter electrode, and an SCE reference electrode. A neutral 0.5 M Na<sub>2</sub>SO<sub>4</sub> electrolyte was used. The light source was a variable-intensity LED, emitting monochromatic 365 nm UV light.

A typical characterization procedure of an electrochemically deposited and annealed  $WO<sub>3</sub>$  or  $WO<sub>3</sub>$ -TiO<sub>2</sub> film consisted of these steps:

Voltammetry experiments were conducted from 0.2 V to 1.6 V (vs. SCE) at 5 mV s<sup>-1</sup>, with a 4-s on/off illumination pulse (2 s on/2 s off). Chronoamperometry at a constant potential of 1.2 V used an illumination pulse of 120 s  $(60 \text{ s on}/60 \text{ s off})$ . Light intensity was set to 30 mW cm<sup>-2</sup>. IMPS was measured at steady state potentials: 0.6 V, 0.8 V, 1.0 V, 1.2 V (100 mW cm−<sup>2</sup> light intensity, 10% amplitude modulation, 10 kHz to 0.1 Hz).

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#### **3. Results and Discussion 3. Results and Discussion**

*3.1. Structure and Surface Morphology 3.1. Structure and Surface Morphology* 

Figure [2](#page-3-0) shows the XRD patterns of electrodeposited WO<sub>3</sub> and WO<sub>3</sub>-TiO<sub>2</sub> films that had been annealed at  $400 °C$  for 2 h. The XRD pattern of the annealed stainless steel substrate is also included. It contains two peaks characteristic of austenite iron (JCPDS strate is also included. It contains two peaks characteristic of austenite iron (JCPDS No. No. 00-023-0298) at 43.5◦ and 50.7◦ , and the substrate peaks also emerge in the diffrac-00-023-0298) at 43.5° and 50.7°, and the substrate peaks also emerge in the diffractograms tograms of the deposited films. The diffractograms of  $WO_3$  and  $WO_3$ -TiO<sub>2</sub> reveal the radically different phase structure of the films.  $WO_3$ -TiO<sub>2</sub> displays all the characteristic peaks of monoclinic WO<sub>3</sub> [\[38](#page-13-10)[,39\]](#page-13-11). However, there are a couple of relatively low-intensity peaks that cannot be attributed to this phase, in particular, at 27.6 $^{\circ}$  and 49.5 $^{\circ}$ . The analysis of WO<sub>3</sub> reveals the origin of these peaks. The diffractogram of WO<sub>3</sub> contains a very intense peak at ~ 24.1°, which may correspond to the (200) face of monoclinic WO<sub>3</sub>. There is also a broad hump at a slightly lower 2Θ of 23.3°, and two small peaks at 27.8° and 49.5°. These reasonably correspond to the unidentified peaks from  $WO_3$ -TiO<sub>2</sub> and have been attributed to hydrated forms of  $WO_3$  in the literature [\[40](#page-13-12)[,41\]](#page-13-13). It is claimed that the formation of WO<sub>3</sub>·xH<sub>2</sub>O can occur through tuning of the intermediate peroxo-complex and thus is related to the  $H_2O_2$  content in the solution [[40\].](#page-13-12) The electrodeposition of a hydrated tungsten oxide has also recently been reported under conditions very similar to those used in this study [[42\]](#page-13-14). Therefore, the formation of monoclinic WO<sub>3</sub> when electrodepositing the composite must have been related to an equilibrium shift in the intermediate peroxo-complex posite must have been related to an equilibrium shift in the intermediate peroxo-complex caused by the interaction of solution species with nanoparticulate TiO<sub>2</sub>, possibly solvation. The crystallite sizes of WO<sub>3</sub>, as approximated by the Scherer equation, are  $41 \pm 9$  nm for WO<sub>3</sub> and 27.4  $\pm$  3.3 nm for WO<sub>3</sub>-TiO<sub>2</sub>. Regarding the presence of TiO<sub>2</sub>, in the diffractogram of the composite film, a very faint hump at ~25◦ can perhaps be distinguished, which could of the composite film, a very faint hump at ~25° can perhaps be distinguished, which could be related to the main (101) peak of anatase  $TiO<sub>2</sub>$ .

<span id="page-3-0"></span>

**Figure 2.** XRD patterns of stainless steel substrate and electrodeposited  $WO_3$  and  $WO_3$ -TiO<sub>2</sub> composite films, after annealing at 400  $^{\circ}$ C.

The visual appearance and SEM surface morphology images of  $WO_3$  and  $WO_3$ -TiO<sub>2</sub> films are shown in Figure [3.](#page-4-0) The  $WO<sub>3</sub>$  film (Figure [3a](#page-4-0)) is inhomogeneously polychromic, which is indicative of a compact thin film of slightly varying thicknesses. The surface morphology, as observed by SEM, is smooth and fine-grained, as was also observed by Pauporté when first reporting on  $WO_3$  electrodeposition by this method [\[34\]](#page-13-7). For both films, signs of emerging cracks can be seen. It was noted earlier [\[43](#page-13-15)[,44\]](#page-13-16) that cracking and

<span id="page-4-0"></span>delamination occur if the layer is thicker (i.e., after longer deposition times). Accordingly, the  $WO_3$ -Ti $O_2$  film is already cracked (Figure 3b), [wh](#page-4-0)ich may indicate a larger thickness. Scattered agglomerates of white particles, which are most likely  $TiO<sub>2</sub>$ , can be observed on the surface. Visually, the film appears "cloudy" because  $TiO<sub>2</sub>$  particles had settled at the surface of the film during deposition. These particles were not always well-adhered to the surface. If the electrodeposition time was prolonged, they would detach from the surface. If the electrodeposition time was prolonged, they would detach from the surface surface when immersed in an electrolyte for further characterization. For this reason, the PEC characterization was carried out on films that had been deposited for 5 min.

 $p$  is the first report on  $W$  electrodeposition by this method  $\tilde{a}$ 



Figure 3. SEM surface morphology images of (a) WO<sub>3</sub>; (b) WO<sub>3</sub>-TiO<sub>2</sub> films electrodeposited for 5 min, (c) SEM image of EDX line scan on a  $WO_3$ -TiO<sub>2</sub> film, (d) EDX line scan data portraying titanium and and tungsten signals. tungsten signals.

The results of an EDX line scan (Figure [3c](#page-4-0),d) show that a  $TiO<sub>2</sub>$  nanoparticle agglomeration was observed on the surface. Titanium was detected, albeit with a very weak signal. It is also interesting to note that the tungsten signal increases along with titanium. This may simply be a topographical effect (i.e., a stronger signal obtained from a point closer to the detector), but it could also indicate that  $WO_3$  had electrodeposited onto TiO<sub>2</sub>.

The near-surface electronic structure of the  $WO_3$  and  $WO_3$ -TiO<sub>2</sub> films was investigated by XPS. It was found that for both films, the core-level structure of W 4f deconvolutes into a single doublet with W  $4f_{7/2}$  and W  $4f_{5/2}$  peaks at binding energies of 35.46 eV and 37.6 eV, respectively (Figure [4a](#page-5-0),b). The position of these peaks, and the spin-orbit splitting energy of 2.14 eV, indicate the total dominance of the  $W^{6+}$  oxidation state [\[45,](#page-13-17)[46\]](#page-13-18). Interestingly, the O 1s core-level spectra reveal that oxygen can be found in at least two oxidation states

<span id="page-5-0"></span>

(Figure [4c](#page-5-0),d). The O 1s peak seen here at 530.22 eV for  $WO_3$  and 530.48 eV for  $WO_3$ -TiO<sub>2</sub> is commonly attributed to lattice oxygen (i.e., the W-O bond) [\[47\]](#page-13-19).

splitting energy of 2.14 eV, indicate the total dominance of the W6+ oxidation state [45,46].

Figure 4. Deconvoluted core-level XPS spectra of W 4f, O 1s, and Ti 2p of a WO<sub>3</sub> film (a,c,e) and WO3-TiO2 film (**b**,**d**,**f**). WO<sup>3</sup> -TiO<sup>2</sup> film (**b**,**d**,**f**).

The origin of the second much broader peak at ~531.4 eV is slightly more complicated to explain. First and foremost, this cannot be Ti-related O because (1) the lattice oxygen peak energy is expected to be around 529–530 eV [\[48\]](#page-13-20) and (2) the same peak also exists for  $WO<sub>3</sub>$  without TiO<sub>2</sub>. A couple of other possibilities that are often reported are oxygen from non-stoichiometric  $WO_{3-x}$  or -OH groups incorporated within the lattice, which can result in peaks in the O 1s spectra at approximately  $532.5 \text{ eV}$  [\[47](#page-13-19)[,49\]](#page-14-0). However, these energies also do not match. A very recent work by T. J. Frankcombe and Y. Liu describes what is

most likely the cause of these peaks [\[50\]](#page-14-1). They found that a peak at this energy (commonly ascribed in the literature to oxygen vacancies) is instead caused by  $O$  1s electrons from water molecules that are chemisorbed to the surface or surface oxygen passivated by hydrogen. This agrees well with the W 4f spectra, which show 100% dominance of  $W^{6+}$  and no trace of reduced oxidation states that would be expected with a large number of oxygen vacancies. The existence of chemisorbed water may also be related to the observation of hydrates is  $V_{\text{S}}$ . in XRD characterization (Figure [2\)](#page-3-0). Lastly, no Ti is found on WO<sub>3</sub>, and only a very weak Ti 2p signal is obtained from the  $WO_3$ -TiO<sub>2</sub> film (Figure [4e](#page-5-0),f). The core level spectrum deconvolutes into one doublet with peaks at  $459.25$  eV and  $465.1$  eV, corresponding to  $Ti<sup>4+</sup>$ from anatase TiO<sub>2</sub> [\[51\]](#page-14-2). An indistinct hump at 461.5 eV may also be observed. It is on<br>The reduction of the Ti<sup>3+</sup> Li<sup>31</sup>. the verge of being noise, but a peak at this position has been attributed to  $Ti^{3+}$  [\[52\]](#page-14-3). The reduction of titanium dioxide nanoparticles is feasible under cathodic conditions in an according in an action of titanium dioxide nanoparticles is feasible under cathodic conditions in an acidic electrolyte, so the signal is attributed to  $Ti^{3+}$  here as well. most likely the cause of these peaks [50]. They found that a peak at this energy (commonly seconvolutes into the vertext which peaks at to  $\frac{1}{2}$  of  $\frac{1$ 

most likely the cause of these peaks [50]. They found that a peak at this energy (commonly ascribed in the literature to oxygen vacancies) is instead caused by O 1s electrons from

# 3.2. Photoelectrochemical Characterization

Initially, the PEC properties of WO<sub>3</sub> and WO<sub>3</sub>-TiO<sub>2</sub> films were characterized by carry-ing out a potential scan with intermittent on/off light pulses (Figure [5a](#page-6-0)). Two differences are immediately apparent: (1) the  $WO_3$ -TiO<sub>2</sub> film generates larger photocurrents throughout the entire potential range and (2) the photocurrent onset potential of the composite is more cathodic (~0.2 V vs. SCE) in comparison to just WO<sub>3</sub> (~0.35 V vs. SCE). As expected, the magnitude of the photocurrent increases with the applied potential as the space-charge layer, where the photoexcitation of charge carriers occurs and becomes wider due to the increasing strength of the applied electrical field. It is also worth noting that both  $WO<sub>3</sub>$ and  $WO_3$ -TiO<sub>2</sub> films show an appreciable amount of surface hole recombination, which is evident from the initial photocurrent overshoot that occurs at the instant light is turned on.

<span id="page-6-0"></span>

Figure 5. Chopped UV LSV measurements at low potentials (a) and high potentials (b) in 0.5 M Na<sub>2</sub>SO<sub>4</sub> of WO<sub>3</sub> and WO<sub>3</sub>-TiO<sub>2</sub> films deposited within 5 min; 5 mV s<sup>-1</sup>, 30 mW cm<sup>-2</sup> light intensity.

An additional difference between the activities of these films is noted at higher applied potentials of  $\geq$  ~1 V vs. SCE, which roughly corresponds to 1.23 V vs. RHE in neutral pH (Figure [5b](#page-6-0)). In this case, particularly for  $WO_3$ -Ti $O_2$ , the electrochemical current rises sharply and is most probably attributable to the onset of electrochemical water splitting/oxygen evolution (OER). A similar effect, although with much lower current densities, is observed for the annealed stainless steel substrate. Thus, the composite films exhibit better electrochemical OER activity. This may simply be an effect of increased surface area [\[53\]](#page-14-4) or the improved OER mechanistic kinetics of  $WO<sub>3</sub>$  [\[54\]](#page-14-5).

Because the profile of the photocurrent pulse provides useful information about the Because the profile of the photocurrent pulse provides useful information about the material photoactivity and surface recombination [\[55\]](#page-14-6), potentiostatic photocurrent pulses material photoactivity and surface recombination [55], potentiostatic photocurrent pulses were obtained at 1.2 V vs. SCE and with a longer on/off cycle duration of 120 s (Figure [6\)](#page-7-0). were obtained at 1.2 V vs. SCE and with a longer on/off cycle duration of 120 s (Figure 6). This made it possible to observe the initial photocurrent overshoot immediately upon This made it possible to observe the initial photocurrent overshoot immediately upon turning on the light, and the subsequent decay over several pulses. turning on the light, and the subsequent decay over several pulses.

<span id="page-7-0"></span>

Figure 6. UV on/off pulses in 0.5 M Na<sub>2</sub>SO<sub>4</sub> of WO<sub>3</sub> and WO<sub>3</sub>-TiO<sub>2</sub> films; 365 nm LED, 30 mW cm<sup>-2</sup> light intensity. light intensity.

At the moment the light is turned on, the photocurrent density  $(j_{t=0})$  is proportional to the maximum amount of photogenerated electron–hole pairs. For a photoanode, the to the maximum amount of photogenerated electron–hole pairs. For a photoanode, the electrons move to the back contact and are registered as photocurrent, and the holes move electrons move to the back contact and are registered as photocurrent, and the holes move toward the electrode–electrolyte interface, where they are either lost to surface recombination or are transferred into the electrolyte to carry out an oxidative photoelectrochemical reaction. After this, the photocurrent decays until a steady state value (j<sub>ss</sub> is assumed to be the photocurrent at the end of the 60 s pulse) is reached. In this case, it is observed that the WO<sub>3</sub>-TiO<sub>2</sub> film displays slightly more relative surface recombination than WO<sub>3</sub>. For example, for the WO<sub>3</sub> film, the j<sub>ss</sub>/j<sub>t=0</sub> ratio is 0.8; whereas, for the comparable composite film, it is lower at 0.77. However, the  $WO_3$ -Ti $O_2$  film also generates larger photocurrents overall (~40 µA cm<sup>-2</sup> compared to just under 12 µA cm<sup>-2</sup> for WO<sub>3</sub>). A similar result was observed by Castro et al., who found that a hydrothermally synthesized WO $_3$  composite with 40 wt% of TiO<sub>2</sub> exhibited larger photocurrents than WO<sub>3</sub> [\[21\]](#page-12-14). Similarly, WO<sub>3</sub>/TiO<sub>2</sub> heterostructures on reduced graphene oxide were shown to generate significantly larger heterostructures on reduced graphene oxide were shown to generate significantly larger photocurrents than the plain films alone [22]. photocurrents than the plain films alone [\[22\]](#page-12-16).

Lastly, similar to the initial jump, there is an "overshoot" to the downside when illumination is turned off. This occurs because of the recombination of holes trapped on the surface with electrons—a phenomenon similar to the decay of the photocurrent.

# *3.3. IMPS Study*

IMPS is a non-stationary photoelectrochemical analysis method, modeled by Ponomarev and Peter [\[56\]](#page-14-7), to determine the features of the space charge region relative to the reaction rate constants. A sinusoidal modulation is applied to the intensity of illumination, and the working electrode under investigation is kept at potentiostatic conditions. In this case, the (photo)current response is proportional entirely to the incident photon flux, i.e., conversion efficiency. It must be noted that here, the light intensity was increased in comparison to previous experiments (30 mW cm<sup>-2</sup> to 100 mW cm<sup>-2</sup>). This is because the signal for IMPS is not obtained from the entire photocurrent, but from the light-perturbation-induced photocurrent, which corresponds to the photocurrent generated

by the amplitude of light intensity modulation (10%). If this light intensity is too low (e.g., 3 mW cm<sup>-2</sup> as 10% of 30 mW cm<sup>-2</sup>), the spectra can be noisy. However, photogener-ated charge carrier recombination is generally larger at higher light intensities [\[57\]](#page-14-8), so care should be taken to find the optimal measurement conditions for a particular system.

IMPS Nyquist spectra obtained for  $WO_3$  are shown in Figure [7,](#page-8-0) where H represents the complex transfer function of the resulting current measured by controlling the light in-the complex transfer function of the resulting current measured by controlling the light tensity, normalized to the geometric electrode area. The first quadrant reflects the response from charge-transfer kinetic and surface recombination. The low-frequency intercept H'<sub>LF</sub> (left side) when H"→0 represents the differential steady-state photocurrent increase due to a differential increase in the light intensity and should be zero if no photoelectrochemical reactions occur.  $\omega$ <sub>LF</sub>, the radial frequency of the highest point of the recombination semicircle, is a product of first-order rate constants of photogenerated charge transfer and recombination and is inversely related to the recombination time constant.

<span id="page-8-0"></span>

Figure 7. Nyquist plots for IMPS of a WO<sub>3</sub> film that had been electrodeposited for 5 min, obtained in  $0.5$  M Na<sub>2</sub>SO<sub>4</sub>, 365 nm LED, 100 mW cm<sup>−2</sup> light intensity with 10% modulation amplitude.

Figure  $7$  for WO<sub>3</sub> shows a non-zero  $H'_{LF}$ , suggesting that an electrochemical reaction occurs at all applied potentials. In Figur[e 8](#page-9-0) for WO3-TiO<sub>2</sub> at 0.6 and 0.8 V, this intercept is nearly zero and thus, electron–hole generation and surface recombination primarily occur. A significant shift in this intercept occurs at 1 V where a significant amount of OER is expected to commence driven by the applied potential. The intercept when the imaginary expected to commence driven by the applied potential. The intercept when the imaginary part is zero at a higher frequency (*right side*) represents the amplitude of the flux of charge part is zero at a higher frequency (*right side*) represents the amplitude of the flux of charge carriers to the surface. This intercept occurs at a value of  $H'$  that can be expressed as Equation (4).

$$
H' = j_h \frac{C_H}{C_H + C_{SC}}
$$
\n
$$
A = \frac{1}{2} C_H + \frac{1}{2} C_H +
$$

where  $j_h$  is the hole current amplitude,  $\textsf{C}_\text{H}$  and  $\textsf{C}_\text{SC}$  are the Helmholtz and space charge layer capacitances, respectively.

Since the Helmholtz layer size is typically smaller than the width of the depletion, the space charge layer, then it is expected that the Helmholtz capacitance is higher,  $C_H > C_{SC}$  [\[58\]](#page-14-9). Then, this intercept equals  $j_h$ , i.e., the hole current. At higher frequencies relative to this intercept, the influence of the capacitance of the semiconductor space charge dominates and the response moves to the lower quadrant, with a negative imaginary transfer function. The discussed parameters that had been obtained from the spectra are presented in Table [1.](#page-9-1) The overall trend is that, as the applied potential is raised,  $H'_{LF}$ increases because more photogenerated holes are transferred as surface recombination is suppressed. Moreover, the IMPS spectra are not normalized to the amplitude of the photogenerated hole current, and therefore, larger IMPS spectra magnitudes of  $WO_3$ -TiO<sub>2</sub>

<span id="page-9-0"></span>

suggest higher conversion efficiencies overall, corresponding well to LSV and potentiostatic pulse data.

Figure 8. Nyquist plots for IMPS of a WO<sub>3</sub>-TiO<sub>2</sub> film that had been electrodeposited for 5 min, obtained in 0.5 M Na<sub>2</sub>SO<sub>4</sub>, 365 nm LED, 100 mW cm<sup>-2</sup> light intensity with 10% modulation amplitude.

E, V	WO <sub>3</sub>			$WO_3$ -TiO <sub>2</sub>		
	$\omega_{LF}$ Rad $s^{-1}$	$H'_{LF}$ $\mu$ A cm <sup>-2</sup> $W^{-1}$ m <sup>2</sup>	Jh, $\mu$ A cm <sup>-2</sup> $W^{-1}$ m <sup>2</sup>	$\omega_{LF}$ Rad $s^{-1}$	$H'_{LF}$ $\mu$ A cm <sup>-2</sup> $W^{-1}$ m <sup>2</sup>	Jh, $\mu$ A cm <sup>-2</sup> $W^{-1}$ m <sup>2</sup>
0.6	69.6	0.058	0.135	546.2	0.004	0.511
0.8	37.0	0.097	0.208	271.3	0.010	0.867
$1.0\,$	26.4	0.180	0.276	162.7	0.284	1.208
1.2	26.4	0.330	0.378	124.0	1.028	1.653

<span id="page-9-1"></span>Table 1. Parameters, obtained from interpretation of IMPS Nyquist plots.

of the electrochemical cell, i.e., the series resistance of the substrate and the total capacitance. **EXECUTE:** WE FORE TREATERT SCHLEME CONSTRUCTED THE CONSTRUCTION CHAPTER KING CHAPTER KING CHAPTER SCHLEME RESERVED BY **EXECUTE** rate constant of charge transfer across this same interface and  $k_{\text{rec}}$ —the rate of electronhole recombination at the semiconductor/electrolyte interface. The spectra shown in Figures 5 and 6 exhibit some characteristic tendencies. Most notably—the radius of the low frequency semicircle decreases as the applied potential is increased. In IMPS, the value of the low frequency intercept  $H'_{LF}$  corresponds to the DC photocurrent, scaled by the 10% modulation, and is expressed as Equation (5). In IMPS spectra, the high-frequency semicircle occurs because of the RC time constant Conversely, the low-frequency semicircle describes the recombination and charge transfer

$$
H'_{LF} = j_h \frac{k_{tr}}{k_{tr} + k_{rec}} \tag{5}
$$

Here,  $k_{tr}/k_{tr} + k_{rec}$  corresponds to  $\eta_{trans}$ —the fraction of hole flux from the space charge region that is involved in the charge transfer reaction, or the transfer efficiency. The transfer efficiency increases with applied potential as photogenerated charge carriers can participate in an electrochemical reaction. Therefore, electron-hole recombination is significant at less positive potentials (the low frequency intercept value is closer to  $0$ ), because there is no electrochemical reaction, but at  $1.2$  V, it decreases considerably with the occurrence of OER.

Several important parameters can be calculated from IMPS spectra. Firstly, the transfer efficiency is an indicator of how many of the holes that had arrived at the semiconductor/electrolyte interface that were transferred into the electrolyte to provide the photoelectrochemical reaction. The data for  $WO_3$  and  $WO_3$ -TiO<sub>2</sub> films are shown in Figure [9.](#page-10-0) As expected, the efficiencies are larger with increasing applied potential. Plain WO<sub>3</sub> films have larger photogenerated transfer efficiency over the entire measured potential range, while the TiO<sub>2</sub>-containing composites have total recombination at lower potentials, but improve rapidly at higher potentials.

<span id="page-10-0"></span>

**Figure 9.** Transfer efficiency values, obtained from IMPS spectra for  $WO_3$  and  $WO_3$ -TiO<sub>2</sub> films.

To fully discern reaction kinetics, the rate constants should be considered. Both *ktr* To fully discern reaction kinetics, the rate constants should be considered. Both *ktr* and *krec* can be calculated from two characteristic points of the IMPS spectrum: the low and *krec* can be calculated from two characteristic points of the IMPS spectrum: the low frequency intercept  $H'_{LF}$  (Equation (5)) and  $\omega_{LF}$ , the radial frequency of the highest point of the recombination semicircle, which is described by Equation (6). of the recombination semicircle, which is described by Equation (6).

$$
\omega_{LF} = k_{tr} + k_{rec} \tag{6}
$$

The spectra were analyzed and the rate constants are presented in Figure [10](#page-11-0) in relation to applied potential. Firstly, it can be seen that for both measured photoelectrodes,  $k_{rec}$ decreases with applied potential (Figure 10a), which is expected from the added potential decreases with applied potential (Figure [10a](#page-11-0)), which is expected from the added potential bias as has been discussed earlier. bias as has been discussed earlier.

The  $k_{rec}$  values of  $WO_3$ -TiO<sub>2</sub> are larger, showing faster recombination kinetics. However, when considering the kinetic transfer constant, it is apparent that  $k_t$  increases significantly for the composite films at higher potentials, compared to  $WO<sub>3</sub>$  (Figure [10b](#page-11-0)). Thus, there is a higher OER reaction rate on the  $WO_3$ -TiO<sub>2</sub> under a biased photoelectrochemical environment. At an applied potential of 1.2 V vs. SCE, the  $k$ <sub>tr</sub> of the composite is three times larger than that of the respective  $WO_3$  film, consistent with the voltammetry result in Figure [5](#page-6-0) and the chronoamperometry in Figure [6.](#page-7-0)

While the addition of discrete  $TiO<sub>2</sub>$  nanoparticles to  $WO<sub>3</sub>$  generated a higher photocurrent compared to  $WO<sub>3</sub>$  alone, it had the negative effect of increasing the recombination rate constant, particularly near open circuit conditions at small applied potentials. The increased recombination rate may be due to the introduction of grain boundaries and local potential perturbation by the particles. The recombination rate constant decreased with applied potential, and in the applied potential range for conventional electrolysis, the  $WO_3$ -TiO<sub>2</sub> composite greatly exceeded the performance of the WO<sub>3</sub>. At an even larger applied poten-tial exceeding 1.2 V vs. SCE, evident in the polarization curve in Figure [5,](#page-6-0) the  $WO_3$ -TiO<sub>2</sub>

composite OER current density greatly exceeded  $WO_3$  even without light; thus, the  $TiO_2$ addition to WO<sub>3</sub> improved its electrocatalytic performance. Georgieva et al. [\[30\]](#page-13-3) also found an enhanced photocatalysis for  $TiO<sub>2</sub>/WO<sub>3</sub>$  layers on stainless steel compared to WO<sub>3</sub> but attributed the enhancement to reduced recombination. The results presented here are in contrast to this reasoning, elucidated by the additional evidence from IMPS. The enhancement in our work is due to larger charge transfer kinetics and not suppressed recombination.

<span id="page-11-0"></span>

Figure 10. Values of (a)  $k_{rec}$  and (b)  $k_{tr}$  for WO<sub>3</sub> and WO<sub>3</sub>-TiO<sub>2</sub> films that had been electrodeposited for 5 min. for 5 min.

# The krec values of WO3-TiO2 are larger, showing faster recombination kinetics. How-**4. Conclusions**

In this study, WO<sub>3</sub> and composite WO<sub>3</sub>-TiO<sub>2</sub> films were prepared by electrodeposition from an acidic peroxo-tungstate precursor solution containing suspended anatase TiO<sub>2</sub> nanoparticles for the formation of the composite. The annealed films were found to be crystalline, with WO<sub>3</sub>-TiO<sub>2</sub> having a distinct monoclinic WO<sub>3</sub> structure, but plain WO<sub>3</sub> displaying characteristics of hydrate formation. XPS showed the dominance of W<sup>6+</sup> in both samples. In addition, the inclusion of  $Ti^{4+}$  and a small amount of  $Ti^{3+}$  was proven by deconvoluting the Ti 2p core level spectra. The films prepared in this study were found to be photoelectrochemically active, and WO<sub>3</sub>-TiO<sub>2</sub> generated significantly larger photocurrent magnitudes compared to the WO<sub>3</sub> alone. However, the WO<sub>3</sub>-TiO<sub>2</sub> composite photocurrents possessed considerable recombination at and near open circuit conditions. While the overall transfer efficiency is higher for WO<sub>3</sub>, applying a bias potential resulted in a significant increase in the OER current density of the  $WO_3$ -TiO<sub>2</sub> composite: it is three times higher at an applied potential of 1.2 V vs. SCE, and an associated expected decrease in the electron– hole recombination rate. The implications of these results show that electrodeposition of a WO<sub>3</sub>-TiO<sub>2</sub> composite from an electrolyte containing a peroxo-tungstate precursor and  $TiO<sub>2</sub>$  nanoparticles result in significant changes in the material's structure, morphology, and consequently, PEC activity. The synthesis of composites is facile, they exhibit improved PEC performance as photoanodes when compared to  $WO_3$ , but also suffer from higher photogenerated electron–hole recombination at lower applied potentials.

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